

Well-Posedness for a New Two-Component Integrable System

Qiong LONG^{1,*}, Chunlai MU¹, Pan ZHENG¹, Shouming ZHOU²

1. College of Mathematics and Statistics, Chongqing University, Chongqing 401331, P. R. China;

2. College of Mathematics, Chongqing Normal University, Chongqing 401331, P. R. China

Abstract In this paper, we consider a new two-component integrable system with cubic nonlinearity, which can be deduced by a curve flow and it is integrable with its Lax pair, bi-Hamiltonian structure, and infinitely many conservation laws. We mainly establish the local well-posedness of this system in a range of the Besov spaces $B_{p,r}^s$ with $s > \max\{2 + \frac{1}{p}, \frac{5}{2}\}$.

Keywords Besov space; two-component integrable system; local well-posedness.

MR(2010) Subject Classification 35A01; 35A02; 35E15; 35L05

1. Introduction

The present paper focuses on the Cauchy problem of the new two-component integrable system with cubic nonlinearity:

$$\begin{cases} m_t = bu_x + \frac{1}{2}[m(uv - u_x v_x)]_x - \frac{1}{2}m(uv_x - u_x v), & t > 0, x \in R \\ n_t = bv_x + \frac{1}{2}[n(uv - u_x v_x)]_x + \frac{1}{2}n(uv_x - u_x v), & t > 0, x \in R \\ m = u - u_{xx}, n = v - v_{xx}, & t > 0, x \in R, \\ u(x, 0) = u_0(x), v(x, 0) = v_0(x) \end{cases} \quad (1)$$

where b is an arbitrary constant. This system was proposed by Qu, Song and Yao in [1]. Note that system (1) can be considered as the two-component case of Eq. (40) in [1]. And Qu et al. [1] showed that Eq. (40) can be deduced by a curve flow $\gamma(x, t)$. Then system (1) can be induced by a curve flow $\gamma(x, t)$, where \vec{m} denotes the curvature vector.

Note that when we take $v = 2$, the system (1) becomes the famous Camassa-Holm equation:

$$m_t + bu_x + 2mu_x + m_x u = 0, \quad (2)$$

which has been derived independently by Fokas and Fuchssteiner [2] by the method of recursion operators, and by Camassa and Holm in [3]. Like the celebrated KdV equation, Camassa-Holm equation describes the unidirectional propagation of shallow water waves over a flat bottom. Here $u(t, x)$ stands for the fluid velocity at time t in the spatial x direction (or equivalently the height of the free surface of water above a flat bottom), b is a constant related to the critical

Received May 16, 2013; Accepted September 11, 2013

Supported by National Natural Science Foundation of China (Grant No. 11371384).

* Corresponding author

E-mail address: longqiong2012@gmail.com (Qiong LONG)

shallow water wave speed. It turns out that it is also a model for the propagation of nonlinear waves in cylindrical hyperplastic rods [4] with $u(x, t)$ representing the radial stretch relative to a pre-stressed state. Since the work of Camassa and Holm, various studies on this equation have been developed. For instance, it has been found that Eq. (2) conforms with many conservation laws, is completely integrable [5] and admits bi-Hamiltonian structures [6], Lax representation, multi-soliton solutions and algebra-geometric solutions [7]. Meanwhile, Eq. (2) possesses smooth solitary wave solutions if $b > 0$ (see [8]) or peakons if $b = 0$ (see [3]). It is also regarded as a model of the geodesic flow for the H^1 right invariant metric on the Bott-Virasoro group if $b > 0$ and on the diffeomorphism group if $b = 0$ (see [9]). The local well-posedness, global existence, blow-up structures and the well-posedness of global weak solutions of Eq. (2) have been given in [10]. The sharpest results for the global existence and blow-up solutions were found in Bressan and Constantin [11].

For $v = 2u$, system (1) is exactly the cubic CH equation:

$$m_t + (u^2 - u_x^2)m_x + 2u_x m^2 + bu_x = 0, \quad (3)$$

which was derived independently by Fuchssteiner [12], Olver and Rosenau [13] and Fokas [14] by applying the general method of tri-Hamiltonian duality to the bi-Hamiltonian representation of the modified Korteweg-de Vries equation. Later, it was derived by Qiao [15] from the two-dimensional Euler equations, where the variables $u(t, x)$ and $m(t, x)$ represent, respectively, the velocity of the fluid and its potential density. In [15, 16], it was shown that equation (3) admits a Lax pair, and hence can be solved by the method of inverse scattering. In [17], the local well-posedness and blow-up scenario to its Cauchy problem were considered. In [18], it was shown that the singularities of the solutions can occur only in the form of wave-breaking, and a new wave-breaking mechanism was described for solutions with certain initial profiles.

If we choose $v = k_1 u + k_2$, system (1) became the generalized CH equation:

$$\begin{cases} m_t + bu_x + \frac{1}{2}k_1[m(u^2 - u_x^2)]_x + \frac{1}{2}k_2(2mu_x + m_x u) = 0, & t > 0, x \in \mathbb{R}, \\ u(x, 0) = u_0(x), & x \in \mathbb{R}, \end{cases} \quad (4)$$

where b, k_1 and k_2 are arbitrary constants, which was proposed by Qiao, Xia and Li [19]. In [19], Qiao et al showed that the above equation was completely integrable, and derived its Lax pair, bi-Hamiltonian structure, peakons, weak kinks, kink-peakon interactional and smooth soliton solutions. Very recently, we have established the well-posedness, blow-up criteria and the lower bound of the maximal time of existence for Eq. (4) in our another paper.

Since the work of CH equation, the study on the two-component CH system has also been remarkably developed [20–22]. In [23], Gui and Liu established the local well-posedness for the two-component Camassa-Holm system. Inspired by the references cited above, the goal of the present paper is to establish the local well-posedness for the strong solutions to the Cauchy problem (1). The proof of the local well-posedness is inspired by the argument of approximate solutions by Danchin [24] in the study of the local well-posedness to the Camassa-Holm equation.

However, one problematic issue is that we here deal with two-component system with a higher order nonlinearity in the Besov spaces, making the proof of several required nonlinear estimates somewhat delicate. These difficulties are nevertheless overcome by carefully estimating each iterative approximation of solutions to (1).

2. Basic facts and notation

For convenience, we need to recall some basic facts on the transport equations and Moser-type estimates.

Definition 2.1 ([25]) (*Besov space*) Let $s \in \mathbb{R}, 1 \leq p, r \leq \infty$. The inhomogeneous Besov space $B_{p,r}^s(\mathbb{R}^d)$ ($B_{p,r}^s$ for short) is defined by

$$B_{p,r}^s \doteq \{f \in \mathcal{S}'(\mathbb{R}^d); \|f\|_{B_{p,r}^s} < \infty\},$$

where

$$\|f\|_{B_{p,r}^s} \doteq \begin{cases} (\sum_{q \in \mathbb{Z}} 2^{qsr} \|\Delta_q f\|_{L^p}^r)^{\frac{1}{r}}, & \text{for } r < \infty, \\ \sup_{q \in \mathbb{Z}} 2^{qs} \|\Delta_q f\|_{L^p}, & \text{for } r = \infty. \end{cases}$$

If $s = \infty, B_{p,r}^\infty \doteq \bigcap_{s \in \mathbb{R}} B_{p,r}^s$.

Definition 2.2 For $T > 0, s \in \mathbb{R}$ and $1 \leq p \leq +\infty$, we set

$$E_{p,r}^s(T) \doteq C([0, T]; B_{p,r}^s) \cap C^1([0, T]; B_{p,r}^{s-1}) \text{ if } r < +\infty,$$

$$E_{p,\infty}^s(T) \doteq L^\infty([0, T]; B_{p,\infty}^s) \cap \text{Lip}([0, T]; B_{p,\infty}^{s-1})$$

and $E_{p,r}^s \doteq \cap_{T>0} E_{p,r}^s(T)$.

Proposition 2.3 ([25, 26]) The following properties hold.

- I) *Density*: if $p, r < \infty$, then $\mathcal{S}(\mathbb{R}^d)$ is dense in $B_{p,r}^s(\mathbb{R}^d)$.
- II) *Sobolev embedding*: if $p_1 \leq p_2$, and $r_1 \leq r_2$, then $B_{p_1,r_1}^s \hookrightarrow B_{p_1,r_2}^{s-d(\frac{1}{p_1}-\frac{1}{p_2})}$. If $s_1 < s_2, 1 \leq p \leq +\infty$ and $1 \leq r_1, r_2 \leq +\infty$, then the embedding $B_{p,r_2}^{s_2} \hookrightarrow B_{p,r_1}^{s_1}$ is locally compact.
- III) *Algebraic properties*: for $s > 0, B_{p,r}^s \cap L^\infty$ is an algebra. Moreover, $B_{p,r}^s$ is an algebra $\Leftrightarrow B_{p,r}^s \hookrightarrow L^\infty \Leftrightarrow s > \frac{d}{p}$ or $(s \geq \frac{d}{p} \text{ and } r = 1)$.
- IV) *Complex interpolation*: if $u \in B_{p,r}^s \cap B_{p,r}^{\tilde{s}}$ and $\theta \in [0, 1], 1 \leq p, r \leq \infty$, then $u \in B_{p,r}^{\theta s + (1-\theta)\tilde{s}}$, and $\|u\|_{B_{p,r}^{\theta s + (1-\theta)\tilde{s}}} \leq \|u\|_{B_{p,r}^s}^\theta \|u\|_{B_{p,r}^{\tilde{s}}}^{1-\theta}$.

Lemma 2.1 ([25, 26]) Suppose that $(p, r) \in [1, \infty]^2$, and $s > -\frac{d}{p}$. Let v be a vector field such that ∇v belongs to $L^1([0, T]; B_{p,r}^{s-1})$ if $s > 1 + \frac{d}{p}$ or to $L^1([0, T]; B_{p,r}^{\frac{d}{p}} \cap L^\infty)$ otherwise. Suppose also that $f_0 \in B_{p,r}^s, F \in L^1([0, T]; B_{p,r}^s)$ and that $f \in L^\infty([0, T]; B_{p,r}^s) \cap C([0, T]; \mathcal{S}')$ solves the d -dimensional linear transport equations

$$\begin{cases} \partial_t f + v \cdot \nabla f = F, \\ f|_{t=0} = f_0. \end{cases} \quad (5)$$

Then there exists a constant C depending only on s, p and d such that the following statements hold:

I) If $r = 1$ or $s \neq 1 + \frac{d}{p}$, then

$$\|f\|_{B_{p,r}^s} \leq \|f_0\|_{B_{p,r}^s} + \int_0^t \|F(\tau)\|_{B_{p,r}^s} d\tau + C \int_0^t V'(\tau) \|f(\tau)\|_{B_{p,r}^s} d\tau,$$

or

$$\|f\|_{B_{p,r}^s} \leq e^{CV(t)} \left(\|f_0\|_{B_{p,r}^s} + \int_0^t e^{-CV(\tau)} \|F(\tau)\|_{B_{p,r}^s} d\tau \right),$$

where $V(t) = \int_0^t \|\nabla v(\tau)\|_{B_{p,r}^{\frac{d}{p}} \cap L^\infty} d\tau$ if $s < 1 + \frac{d}{p}$ and $V(t) = \int_0^t \|\nabla v(\tau)\|_{B_{p,r}^{s-1}} d\tau$ else.

II) If $s \leq 1 + \frac{d}{p}$ and, in addition, $\nabla f_0 \in L^\infty$, $\nabla f \in L^\infty([0, T] \times \mathbb{R}^d)$ and $\nabla F \in L^1([0, T]; L^\infty)$, then

$$\|f\|_{B_{p,r}^s} + \|\nabla f(t)\|_{L^\infty} \leq e^{CV(t)} \left(\|f_0\|_{B_{p,r}^s} + \|\nabla f_0\|_{L^\infty} + \int_0^t e^{-CV(\tau)} \|F(\tau)\|_{B_{p,r}^s} + \|\nabla F(\tau)\|_{L^\infty} d\tau \right)$$

with $V(t) = \int_0^t \|\nabla v(\tau)\|_{B_{p,r}^{\frac{d}{p}} \cap L^\infty} d\tau$.

III) If $f = v$, then for all $s > 0$, the estimate (6) holds with $V(t) = \int_0^t \|\partial_x u(\tau)\|_{L^\infty} d\tau$.

IV) If $r < +\infty$, then $f \in C([0, T]; B_{p,r}^s)$. If $r = +\infty$, then $f \in C([0, T]; B_{p,1}^{s'})$ for all $s' < s$.

Lemma 2.2 ([26]) Let $(p, p_1, r) \in [1, +\infty]^3$. Assume that $s > -d \min\{\frac{1}{p}, \frac{1}{p'}\}$ with $p' := (1 - \frac{1}{p})^{-1}$. Let $f_0 \in B_{p,r}^s$ and $F \in L^1([0, T]; B_{p,r}^s)$. Let v be a time dependent vector field such that $v \in L^\rho([0, T]; B_{\infty,\infty}^{-M})$ for some $\rho > 1, M > 0$, and $\nabla v \in L^1([0, T]; B_{p_1,\infty}^{\frac{d}{p_1}} \cap L^\infty)$ if $s < 1 + \frac{d}{p_1}$, and $\nabla v \in L^1([0, T]; B_{p_1,r}^{s-1})$ if $s > 1 + \frac{d}{p_1}$ or $s = 1 + \frac{d}{p_1}$ and $r = 1$. Then the transport equations (5) has a unique solution $f \in L^\infty([0, T]; B_{p,r}^s) \cap (\cap_{s' < s} C([0, T]; B_{p,1}^{s'}))$ and the inequalities in Lemma 2.1 hold true. If, moreover, $r < \infty$, then we have $f \in C([0, T]; B_{p,r}^s)$.

Lemma 2.3 ([13]) (1-D Moser-type estimates) Assume that $1 \leq p, r \leq +\infty$. Then the following estimates hold:

(I) For $s > 0$, $\|fg\|_{B_{p,r}^s} \leq C(\|f\|_{B_{p,r}^s} \|g\|_{L^\infty} + \|g\|_{B_{p,r}^s} \|f\|_{L^\infty});$

(II) For $s_1 \leq \frac{1}{p}, s_2 > \frac{1}{p}$ ($s_2 \geq \frac{1}{p}$ if $r = 1$) and $s_1 + s_2 > 0$,

$$\|fg\|_{B_{p,r}^{s_1}} \leq C \|f\|_{B_{p,r}^{s_1}} \|g\|_{B_{p,r}^{s_2}},$$

where the constant C is independent of f and g .

3. Local well-posedness

In this section, we will prove the local well-posedness of the Cauchy problem (1). For

convenience, applying the operator $(1 - \partial_x^2)^{-1}$ to both sides of the system (1), we have

$$\begin{cases} u_t = (1 - \partial_x^2)^{-1}(bu_x) + \partial_x(1 - \partial_x^2)^{-1}\left[u^2 + uu_x + \frac{1}{2}u_x^2\right]v + \\ \quad (1 - \partial_x^2)^{-1}\left(\frac{1}{3}u_x^2v - u^2v\right) + \left(uv - \frac{1}{3}u_xv_x\right)u_x, \\ v_t = (1 - \partial_x^2)^{-1}(bv_x) + \partial_x(1 - \partial_x^2)^{-1}\left[v^2 + vv_x + \frac{1}{2}v_x^2\right]u + \\ \quad (1 - \partial_x^2)^{-1}\left(\frac{1}{3}uv_x^2 - uv^2\right) + \left(uv - \frac{1}{3}u_xv_x\right)v_x, \end{cases} \quad (7)$$

which enables us to define the solutions of the Cauchy problem (1).

Theorem 3.1 Suppose that $1 \leq p, r \leq +\infty$, $s > \max\{2 + \frac{1}{p}, \frac{5}{2}\}$ and $(u_0, v_0) \in B_{p,r}^s \times B_{p,r}^s$. Then there exists a time $T > 0$ such that the initial-value problem (1.1) has a unique solution $(u, v) \in E_{p,r}^s(T) \times E_{p,r}^s(T)$, and the map $(u_0, v_0) \mapsto (u, v)$ is continuous from a neighborhood of $(u_0, v_0) \in B_{p,r}^s \times B_{p,r}^s$ into

$$C([0, T]; B_{p,r}^{s'}) \cap C^1([0, T]; B_{p,r}^{s'-1}) \times C([0, T]; B_{p,r}^{s'}) \cap C^1([0, T]; B_{p,r}^{s'-1})$$

for every $s' < s$ when $r = +\infty$ and $s' = s$ whereas $r < +\infty$.

To prove this theorem, we need to prove the following Proposition and Lemma first. Uniqueness and continuity with respect to the initial data are an immediate consequence of the following result.

Proposition 3.2 Let $1 \leq p, r \leq +\infty$ and $s > \max\{2 + \frac{1}{p}, \frac{5}{2}\}$. Let $(u^{(1)}, v^{(1)}), (u^{(2)}, v^{(2)})$ be two given solutions of the initial-value problem (1) with the initial data $(u_0^{(1)}, v_0^{(1)}), (u_0^{(2)}, v_0^{(2)}) \in B_{p,r}^s \times B_{p,r}^s$ satisfying $u^{(1)}, v^{(1)}, u^{(2)}, v^{(2)} \in L^\infty([0, T]; B_{p,r}^s) \cap C([0, T]; \mathcal{S}')$. Then for every $t \in [0, T]$:

$$\begin{aligned} & \| (u^{(1)} - u^{(2)})(t) \|_{B_{p,r}^{s-1}} + \| (v^{(1)} - v^{(2)})(t) \|_{B_{p,r}^{s-1}} \\ & \leq (\| (u_0^{(1)} - u_0^{(2)}) \|_{B_{p,r}^{s-1}} + \| (v_0^{(1)} - v_0^{(2)}) \|_{B_{p,r}^{s-1}}) \cdot \exp \left\{ C \int_0^t (\| u^{(1)}(\tau) \|_{B_{p,r}^s}^2 + \| u^{(2)}(\tau) \|_{B_{p,r}^s}^2 + \right. \\ & \quad \left. \| v^{(1)}(\tau) \|_{B_{p,r}^s}^2 + \| v^{(2)}(\tau) \|_{B_{p,r}^s}^2 + |b|) d\tau \right\}, \end{aligned} \quad (8)$$

where $C > 0$ is a generic constant only depending on p, r, s .

Proof Denote $u^{(12)} := u^{(2)} - u^{(1)}$ and $v^{(12)} := v^{(2)} - v^{(1)}$. It is obvious that

$$u^{(12)}, v^{(12)} \in L^\infty([0, T]; B_{p,r}^s) \cap C([0, T]; \mathcal{S}'),$$

which along with the equivalent formulation (3.1) of (1.1) implies that

$$(u^{(12)}, v^{(12)}) \in C([0, T]; B_{p,r}^{s-1})$$

and $(u^{(12)}, v^{(12)})$ solves the transport equation

$$\begin{cases} \partial_t u^{(12)} + [\frac{1}{3}(u_x^{(1)} + u_x^{(2)})v_x^{(1)} - u^{(1)}v^{(2)}]\partial_x u^{(12)} = f(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)}), \\ \partial_t v^{(12)} + [\frac{1}{3}(v_x^{(1)} + v_x^{(2)})u_x^{(1)} - v^{(1)}u^{(2)}]\partial_x v^{(12)} = g(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)}), \\ u^{(12)}|_{t=0} = u_0^{(12)} := u_0^{(2)} - u_0^{(1)}, \\ v^{(12)}|_{t=0} = v_0^{(12)} := v_0^{(2)} - v_0^{(1)} \end{cases} \quad (9)$$

with

$$\begin{aligned} f(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)}) &= \partial_x(1 - \partial_x^2)^{-1} \left[bu^{(12)} + (u^{(1)} + u^{(2)} + u_x^{(2)})v^{(2)}u^{(12)} + v^{(12)} \times \right. \\ &\quad \left. \left(\frac{1}{2}(u_x^{(2)})^2 + (u^{(1)})^2 + u^{(1)}u_x^{(1)} \right) + (u^{(1)}v^{(2)} + \frac{1}{2}(u_x^{(1)} + u_x^{(2)})v^{(1)})\partial_x u^{(12)} \right] + \\ &\quad (1 - \partial_x^2)^{-1} \left[\frac{1}{3}(u_x^{(1)} + u_x^{(2)})v^{(1)}\partial_x u^{(12)} + \left(\frac{1}{3}(u_x^{(2)})^2 - (u^{(1)})^2 \right)v^{(12)} - \right. \\ &\quad \left. (u^{(1)} + u^{(2)})v^{(2)}u^{(12)} \right] + u_x^{(2)}u^{(12)}v^{(2)} + u^{(1)}u_x^{(1)}v^{(12)} - \frac{1}{3}(u_x^{(2)})^2v_x^{(12)}, \\ g(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)}) &= \partial_x(1 - \partial_x^2)^{-1} \left[bv^{(12)} + (v^{(1)} + v^{(2)} + v_x^{(2)})u^{(2)}v^{(12)} + u^{(12)} \times \right. \\ &\quad \left. \left(\frac{1}{2}(v_x^{(2)})^2 + (v^{(1)})^2 + v^{(1)}v_x^{(1)} \right) + (v^{(1)}u^{(2)} + \frac{1}{2}(v_x^{(1)} + v_x^{(2)})u^{(1)})\partial_x v^{(12)} \right] + \\ &\quad (1 - \partial_x^2)^{-1} \left[\frac{1}{3}(v_x^{(1)} + v_x^{(2)})u^{(1)}\partial_x v^{(12)} + \left(\frac{1}{3}(v_x^{(2)})^2 - (v^{(1)})^2 \right)u^{(12)} - \right. \\ &\quad \left. (v^{(1)} + v^{(2)})u^{(2)}v^{(12)} \right] + v_x^{(2)}v^{(12)}u^{(2)} + v^{(1)}v_x^{(1)}u^{(12)} - \frac{1}{3}(v_x^{(2)})^2u_x^{(12)}. \end{aligned}$$

Thanks to the transport theory in Lemma 2.1, we get

$$\begin{aligned} \|u^{(12)}(t)\|_{B_{p,r}^{s-1}} &\leq \|u_0^{(12)}\|_{B_{p,r}^{s-1}} + \int_0^t \|f(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)})\|_{B_{p,r}^{s-1}} d\tau + \\ &\quad C \int_0^t \left\| \frac{1}{3}(u_x^{(1)} + u_x^{(2)})v_x^{(1)} - u^{(1)}v^{(2)} \right\|_{B_{p,r}^{s-1}} \cdot \|u^{(12)}(\tau)\|_{B_{p,r}^{s-1}} d\tau. \end{aligned} \quad (10)$$

$$\begin{aligned} \|v^{(12)}(t)\|_{B_{p,r}^{s-1}} &\leq \|v_0^{(12)}\|_{B_{p,r}^{s-1}} + \int_0^t \|g(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)})\|_{B_{p,r}^{s-1}} d\tau + \\ &\quad C \int_0^t \left\| \frac{1}{3}(v_x^{(1)} + v_x^{(2)})u_x^{(1)} - v^{(1)}u^{(2)} \right\|_{B_{p,r}^{s-1}} \cdot \|v^{(12)}(\tau)\|_{B_{p,r}^{s-1}} d\tau. \end{aligned} \quad (11)$$

Applying the product law in the Besov spaces, we have

$$\left\| \frac{1}{3}(u_x^{(1)} + u_x^{(2)})v_x^{(1)} - u^{(1)}v^{(2)} \right\|_{B_{p,r}^{s-1}} \leq C(\|u^{(1)}\|_{B_{p,r}^s}^2 + \|u^{(2)}\|_{B_{p,r}^s}^2 + \|v^{(1)}\|_{B_{p,r}^s}^2). \quad (12)$$

$$\begin{aligned} &\|u_x^{(2)}u^{(12)}v^{(2)} + u^{(1)}u_x^{(1)}v^{(12)} - \frac{1}{3}(u_x^{(2)})^2v_x^{(12)}\|_{B_{p,r}^{s-1}} \\ &\leq C(\|u^{(2)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(2)}\|_{B_{p,r}^{s-1}}^2)\|u^{(12)}\|_{B_{p,r}^{s-1}} + C\|u^{(2)}\|_{B_{p,r}^s}^2\|v^{(12)}\|_{B_{p,r}^{s-1}}, \end{aligned}$$

$$\begin{aligned} &\left\| (1 - \partial_x^2)^{-1} \left[\frac{1}{3}(u_x^{(1)} + u_x^{(2)})v^{(1)}\partial_x u^{(12)} + \left(\frac{1}{3}(u_x^{(2)})^2 - (u^{(1)})^2 \right)v^{(12)} - (u^{(1)} + u^{(2)})v^{(2)}u^{(12)} \right] \right\|_{B_{p,r}^{s-1}} \\ &\leq C \left[(\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(1)}\|_{B_{p,r}^{s-1}}^2)\|u^{(12)}\|_{B_{p,r}^{s-1}} + (\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^{s-1}}^2)\|v^{(12)}\|_{B_{p,r}^{s-1}} \right], \end{aligned}$$

and

$$\begin{aligned} & \left\| \partial_x (1 - \partial_x^2)^{-1} [bu^{(12)} + (u^{(1)} + u^{(2)} + u_x^{(2)})v^{(2)}u^{(12)} + v^{(12)} \cdot (\frac{1}{2}(u_x^{(2)})^2 + (u^{(1)})^2 + u^{(1)}u_x^{(1)}) + \right. \\ & \quad \left. (u^{(1)}v^{(2)} + \frac{1}{2}(u_x^{(1)} + u_x^{(2)})v^{(1)})\partial_x u^{(12)}] \right\|_{B_{p,r}^{s-1}} \\ & \leq C(\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(2)}\|_{B_{p,r}^{s-1}}^2 + |b|)\|u^{(12)}\|_{B_{p,r}^{s-1}} + \\ & \quad C(\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^{s-1}}^2)\|v^{(12)}\|_{B_{p,r}^{s-1}}, \end{aligned}$$

which leads to

$$\begin{aligned} \|f(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)})\|_{B_{p,r}^{s-1}} & \leq C \left[(\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^s}^2 + \|v^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(2)}\|_{B_{p,r}^s}^2 + \right. \\ & \quad \left. |b|)\|u^{(12)}\|_{B_{p,r}^{s-1}} + (\|u^{(1)}\|_{B_{p,r}^s}^2 + \|u^{(2)}\|_{B_{p,r}^s}^2)\|v^{(12)}\|_{B_{p,r}^{s-1}} \right]. \quad (13) \end{aligned}$$

Similarly, we have

$$\begin{aligned} \|g(u^{(1)}, u^{(2)}, v^{(1)}, v^{(2)})\|_{B_{p,r}^{s-1}} & \leq C \left[(\|u^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|u^{(2)}\|_{B_{p,r}^s}^2 + \|v^{(1)}\|_{B_{p,r}^{s-1}}^2 + \|v^{(2)}\|_{B_{p,r}^s}^2 + \right. \\ & \quad \left. |b|)\|v^{(12)}\|_{B_{p,r}^{s-1}} + (\|v^{(1)}\|_{B_{p,r}^s}^2 + \|v^{(2)}\|_{B_{p,r}^s}^2)\|u^{(12)}\|_{B_{p,r}^{s-1}} \right]. \quad (14) \end{aligned}$$

Then, inserting Eq. (12–14) into (10) and (11) implies that

$$\begin{aligned} \|u^{(12)}(t)\|_{B_{p,r}^{s-1}} + \|v^{(12)}(t)\|_{B_{p,r}^{s-1}} & \leq (\|u^{(12)}(0)\|_{B_{p,r}^{s-1}} + \|v^{(12)}(0)\|_{B_{p,r}^{s-1}}) + \\ & \quad C \int_0^t (\|u^{(1)}(\tau)\|_{B_{p,r}^s}^2 + \|u^{(2)}(\tau)\|_{B_{p,r}^s}^2 + \|v^{(1)}(\tau)\|_{B_{p,r}^s}^2 + \\ & \quad \|v^{(2)}(\tau)\|_{B_{p,r}^s}^2 + |b|) \cdot (\|u^{(12)}(t)\|_{B_{p,r}^{s-1}} + \|v^{(12)}(t)\|_{B_{p,r}^{s-1}}) d\tau. \end{aligned}$$

Applying Gronwall's inequality, we get (8).

Next, we use the classical Friedrichs regularization method to construct the approximate solutions to the Cauchy problem (1).

Lemma 3.3 *Let $1 \leq p, r \leq +\infty, s > \max\{2 + \frac{1}{p}, \frac{5}{2}\}, (u_0, v_0) \in B_{p,r}^s \times B_{p,r}^s$. Assume that $u^{(0)} = v^{(0)} = 0$. There exists a sequence of smooth functions $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}} \in C(\mathbb{R}^+; B_{p,r}^\infty)^2$ solving the following linear transport equations by induction:*

$$\begin{cases} \partial_t m^{(n+1)} + \frac{1}{2}(u_x^{(n)}v_x^{(n)} - u^{(n)}v^{(n)})\partial_x m^{(n+1)} = P(u^{(n)}, v^{(n)}), & t > 0, x \in \mathbb{R}, \\ \partial_t n^{(n+1)} + \frac{1}{2}(u_x^{(n)}v_x^{(n)} - u^{(n)}v^{(n)})\partial_x n^{(n+1)} = Q(u^{(n)}, v^{(n)}), & t > 0, x \in \mathbb{R}, \\ u^{(n+1)}|_{t=0} = u_0^{(n+1)}(x) = S_{n+1}u_0, & x \in \mathbb{R}, \\ v^{(n+1)}|_{t=0} = v_0^{(n+1)}(x) = S_{n+1}v_0, & x \in \mathbb{R}, \end{cases} \quad (15)$$

where

$$\begin{cases} P(u^{(n)}, v^{(n)}) = bu_x^{(n)} + \frac{1}{2}(2u_x^{(n)}v^{(n)} - (u_x^{(n)}v_x^{(n)})_x)m^{(n)}, \\ Q(u^{(n)}, v^{(n)}) = bv_x^{(n)} + \frac{1}{2}(2v_x^{(n)}u^{(n)} - (u_x^{(n)}v_x^{(n)})_x)n^{(n)}. \end{cases} \quad (16)$$

Moreover, there exists a $T > 0$ such that the solutions satisfy the following properties:

- (I) $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is uniformly bounded in $E_{p,r}^s(T) \times E_{p,r}^s(T)$.
 (II) $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is a Cauchy sequence in $C([0, T]; B_{p,r}^\infty)^2$.

Proof As all data $(S_{n+1}u_0, S_{n+1}v_0)$ belongs to $B_{p,r}^\infty \times B_{p,r}^\infty$, Lemma 2.2 enables us to show by induction that for all $n \in \mathbb{N}$, the equation (15) has a global solution which belongs to $C(\mathbb{R}; B_{p,r}^\infty)^2$. Applying Lemma 2.1 to (15), we get for all $n \in \mathbb{N}$:

$$\begin{aligned} & \|m^{(n+1)}(t)\|_{B_{p,r}^{s-2}} + \|n^{(n+1)}(t)\|_{B_{p,r}^{s-2}} \\ & \leq \exp\left(C \int_0^t \frac{1}{2} \|u_x^{(n)} v_x^{(n)} - u^{(n)} v^{(n)}\|_{B_{p,r}^{s-1}} d\tau\right) \cdot (\|S_{n+1}u_0\|_{B_{p,r}^s} + \|S_{n+1}v_0\|_{B_{p,r}^s}) + \\ & \quad C \int_0^t \exp\left(C \int_\tau^t \frac{1}{2} \|u_x^{(n)} v_x^{(n)} - u^{(n)} v^{(n)}\|_{B_{p,r}^{s-1}} d\tau'\right) \cdot \\ & \quad (\|P(u^{(n)}, v^{(n)})\|_{B_{p,r}^{s-2}} + \|Q(u^{(n)}, v^{(n)})\|_{B_{p,r}^{s-2}}) d\tau. \end{aligned} \quad (17)$$

Thanks to the product law in Besov spaces, one has

$$\begin{aligned} & \|u_x^{(n)} v_x^{(n)} - u^{(n)} v^{(n)}\|_{B_{p,r}^{s-1}} \leq C(\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s})^2, \\ & \|P(u^{(n)}, v^{(n)})\|_{B_{p,r}^{s-2}} + \|Q(u^{(n)}, v^{(n)})\|_{B_{p,r}^{s-2}} \\ & \leq C[(\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s}) + (\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s})^2 + (\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s})^4], \end{aligned}$$

which combined with (17) leads to

$$\begin{aligned} & \|u^{(n+1)}(t)\|_{B_{p,r}^s} + \|v^{(n+1)}(t)\|_{B_{p,r}^s} \\ & \leq \exp\left(C \int_0^t (\|u^{(n)}(\tau)\|_{B_{p,r}^s} + \|v^{(n)}(\tau)\|_{B_{p,r}^s})^2 d\tau\right) \cdot (\|u_0\|_{B_{p,r}^s} + \|v_0\|_{B_{p,r}^s}) + \\ & \quad C \int_0^t \exp\left(C \int_\tau^t (\|u^{(n)}(\tau')\|_{B_{p,r}^s} + \|v^{(n)}(\tau')\|_{B_{p,r}^s})^2 d\tau'\right) \cdot [(\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s}) + \\ & \quad (\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s})^2 + (\|u^{(n)}\|_{B_{p,r}^s} + \|v^{(n)}\|_{B_{p,r}^s})^4] d\tau. \end{aligned} \quad (18)$$

Let us choose a $T > 0$ such that

$$T \leq \min\left\{\frac{3-2M_0}{4C(1+6M_0)}, \frac{1}{16CM_0^2}\right\},$$

where $M_0 = \|u_0\|_{B_{p,r}^s} + \|v_0\|_{B_{p,r}^s}$ and suppose by induction that for all $t \in [0, T]$

$$\|u^{(n)}(t)\|_{B_{p,r}^s} + \|v^{(n)}(t)\|_{B_{p,r}^s} \leq \frac{2M_0}{(1-16CM_0^2t)^{\frac{1}{2}}}. \quad (19)$$

Then we obtain from (19) that for any $0 \leq \tau \leq t$

$$C \int_\tau^t (\|u^{(n)}(\tau')\|_{B_{p,r}^s} + \|v^{(n)}(\tau')\|_{B_{p,r}^s})^2 d\tau' = \frac{1}{4} [\ln(1-16CM_0^2\tau) - \ln(1-16CM_0^2t)].$$

And inserting the above inequality and (19) into (18) leads to

$$\begin{aligned} & \|u^{(n+1)}(t)\|_{B_{p,r}^s} + \|v^{(n+1)}(t)\|_{B_{p,r}^s} \\ & \leq (1-16CM_0^2t)^{-\frac{1}{4}} \left[M_0 + C \int_0^t (2M_0(1-16CM_0^2t)^{-\frac{1}{4}} + 4M_0^2(1-16CM_0^2t)^{-\frac{3}{4}} + \right. \end{aligned}$$

$$\begin{aligned} & 16M_0^4(1 - 16CM_0^2t)^{-\frac{7}{4}}d\tau] \\ & \leq \frac{8C(1 + 6M_0)M_0t + 4M_0^2}{3(1 - 16CM_0^2t)^{\frac{1}{2}}}, \end{aligned}$$

which implies

$$\|u^{(n+1)}(t)\|_{B_{p,r}^s} + \|v^{(n+1)}(t)\|_{B_{p,r}^s} \leq \frac{2M_0}{(1 - 16CM_0^2t)^{\frac{1}{2}}}.$$

Therefore, $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is uniformly bounded in $C([0, T]; B_{p,r}^s)^2$.

Since $B_{p,r}^s$ is an algebra, one can deduce from the Moser-type estimates (see Lemma 2.3 (II)) and equation (15) that

$$(\partial_t u^{(n+1)}, \partial_t v^{(n+1)}) \in C([0, T]; B_{p,r}^{s-1})^2$$

uniformly bounded, which yields that the sequence $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is uniformly bounded in $E_{p,r}^s(T) \times E_{p,r}^s(T)$.

Now, let us prove that $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is a Cauchy sequence in $C([0, T]; B_{p,r}^{s-1})^2$. By (15), we have that, for all $n, k \in \mathbb{N}$,

$$\begin{cases} [\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x](m^{(n+k+1)} - m^{(n+1)}) = F(u^{(n+k)}, u^{(n)}, v^{(n+k)}, v^{(n)}), \\ [\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x](n^{(n+k+1)} - n^{(n+1)}) = G(u^{(n+k)}, u^{(n)}, v^{(n+k)}, v^{(n)}), \end{cases} \quad (20)$$

where

$$\begin{aligned} & F(u^{(n+k)}, u^{(n)}, v^{(n+k)}, v^{(n)}) \\ &= b(u_x^{(n+k)} - u_x^{(n)}) + \frac{1}{2}[(u_x^{(n+k)}v_x^{(n+k)} - u_x^{(n)}v_x^{(n)}) - (u^{(n+k)}v^{(n+k)} - u^{(n)}v^{(n)})]\partial_x m^{(n+1)} + \\ & \quad \frac{1}{2}[2u_x^{(n)}v^{(n)} - (u_x^{(n)}v_x^{(n)})_x](m^{(n+k)} - m^{(n)}) + \frac{1}{2}[2[(u_x^{(n+k)} - u_x^{(n)})v^{(n+k)} + (v^{(n+k)} - v^{(n)})u_x^{(n)}] - \\ & \quad [(u_x^{(n+k)} - u_x^{(n)})v_x^{(n+k)} + (v_x^{(n+k)} - v_x^{(n)})u_x^{(n)}]_x]m^{(n+k)}, \\ & G(u^{(n+k)}, u^{(n)}, v^{(n+k)}, v^{(n)}) \\ &= b(v_x^{(n+k)} - v_x^{(n)}) + \frac{1}{2}[(u_x^{(n+k)}v_x^{(n+k)} - u_x^{(n)}v_x^{(n)}) - (u^{(n+k)}v^{(n+k)} - u^{(n)}v^{(n)})]\partial_x n^{(n+1)} + \\ & \quad \frac{1}{2}[2v_x^{(n)}u^{(n)} - (u_x^{(n)}v_x^{(n)})_x](n^{(n+k)} - n^{(n)}) + \frac{1}{2}[2[(v_x^{(n+k)} - v_x^{(n)})u^{(n+k)} + (u^{(n+k)} - u^{(n)})v_x^{(n)}] - \\ & \quad [(v_x^{(n+k)} - v_x^{(n)})u_x^{(n+k)} + (u_x^{(n+k)} - u_x^{(n)})v_x^{(n)}]_x]n^{(n+k)}. \end{aligned}$$

Note that the equation (20) is equivalent to the following one:

$$\begin{cases} (1 - \partial_x^2)[[\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x](u^{(n+k+1)} - u^{(n+1)})] = \overline{F}^{(n,k)} \\ (1 - \partial_x^2)[[\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x](v^{(n+k+1)} - v^{(n+1)})] = \overline{G}^{(n,k)} \end{cases} \quad (21)$$

with

$$\overline{F}^{(n,k)} = (1 - \partial_x^2)[\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(u^{(n+k+1)} - u^{(n+1)})] -$$

$$\begin{aligned} & \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(1 - \partial_x^2)(u^{(n+k+1)} - u^{(n+1)}) + F, \\ \overline{G}^{(n,k)} = & (1 - \partial_x^2)\left[\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(v^{(n+k+1)} - v^{(n+1)})\right] - \\ & \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(1 - \partial_x^2)(v^{(n+k+1)} - v^{(n+1)}) + G. \end{aligned}$$

Applying the operator $(1 - \partial_x^2)^{-1}$ to (21), one obtains that

$$\begin{cases} \left[\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x\right](u^{(n+k+1)} - u^{(n+1)}) = (1 - \partial_x^2)^{-1}\overline{F}^{(n,k)} \\ \left[\partial_t + \frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x\right](v^{(n+k+1)} - v^{(n+1)}) = (1 - \partial_x^2)^{-1}\overline{G}^{(n,k)}. \end{cases} \quad (22)$$

Thanks to Lemma 2.1, for every $t \in [0, T]$, one gets

$$\begin{aligned} & \exp\left(-C \int_0^t \left\|\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\right\|_{B_{p,r}^{s-1}} d\tau\right) \cdot \|(u^{(n+k+1)} - u^{(n+1)})(t)\|_{B_{p,r}^{s-1}} \\ & \leq C \int_0^t \exp\left(-C \int_0^\tau \left\|\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\right\|_{B_{p,r}^{s-1}} d\tau'\right) \cdot \|\overline{F}^{(n,k)}\|_{B_{p,r}^{s-3}} d\tau + \\ & \quad \|u_0^{(n+k+1)} - u_0^{(n+1)}\|_{B_{p,r}^{s-1}}, \\ & \exp\left(-C \int_0^t \left\|\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\right\|_{B_{p,r}^{s-1}} d\tau\right) \cdot \|(v^{(n+k+1)} - v^{(n+1)})(t)\|_{B_{p,r}^{s-1}} \\ & \leq C \int_0^t \exp\left(-C \int_0^\tau \left\|\frac{1}{2}(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\right\|_{B_{p,r}^{s-1}} d\tau'\right) \cdot \|\overline{G}^{(n,k)}\|_{B_{p,r}^{s-3}} d\tau + \\ & \quad \|v_0^{(n+k+1)} - v_0^{(n+1)}\|_{B_{p,r}^{s-1}}. \end{aligned} \quad (23)$$

Since $s > \max\{2 + \frac{1}{p}, \frac{5}{2}\}$, by the product law in Besov spaces, one has

$$\begin{aligned} & \|(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(u^{(n+k+1)} - u^{(n+1)})\|_{B_{p,r}^{s-1}} \\ & \leq C\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}}(\|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2), \\ & \|(u_x^{(n+k)}v_x^{(n+k)} - u^{(n+k)}v^{(n+k)})\partial_x(1 - \partial_x^2)(u^{(n+k+1)} - u^{(n+1)})\|_{B_{p,r}^{s-3}} \\ & \leq C\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}}(\|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2), \end{aligned}$$

and

$$\begin{aligned} \|F\|_{B_{p,r}^{s-3}} & \leq C\|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}}(\|u^{(n)}\|_{B_{p,r}^s}^2 + \|u^{(n+1)}\|_{B_{p,r}^s}^2 + \|u^{(n+k)}\|_{B_{p,r}^s}^2 + \\ & \quad \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n)}\|_{B_{p,r}^s}^2 + |b|) + C\|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}}(\|u^{(n)}\|_{B_{p,r}^s}^2 + \\ & \quad \|u^{(n+1)}\|_{B_{p,r}^s}^2 + \|u^{(n+k)}\|_{B_{p,r}^s}^2), \\ \|G\|_{B_{p,r}^{s-3}} & \leq C\|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}}(\|v^{(n)}\|_{B_{p,r}^s}^2 + \|v^{(n+1)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \\ & \quad \|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|u^{(n)}\|_{B_{p,r}^s}^2 + |b|) + C\|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}}(\|v^{(n)}\|_{B_{p,r}^s}^2 + \\ & \quad \|v^{(n+1)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2). \end{aligned}$$

Hence, we obtain that

$$\begin{aligned} & \|\overline{F}^{(n,k)}\|_{B_{p,r}^{s-3}} + \|\overline{G}^{(n,k)}\|_{B_{p,r}^{s-3}} \\ & \leq C[\|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}} + \|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}}](\|u^{(n)}\|_{B_{p,r}^s}^2 + \|u^{(n+1)}\|_{B_{p,r}^s}^2 + \|u^{(n+k)}\|_{B_{p,r}^s}^2 + \\ & \quad \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n)}\|_{B_{p,r}^s}^2 + |b|) + C[\|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}} + \|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}}](\|v^{(n)}\|_{B_{p,r}^s}^2 + \\ & \quad \|v^{(n+1)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|u^{(n)}\|_{B_{p,r}^s}^2 + |b|). \end{aligned}$$

$$\begin{aligned} & \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+1)}\|_{B_{p,r}^s}^2 + \|v^{(n)}\|_{B_{p,r}^s}^2 + |b| + (\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}} + \\ & \|v^{(n+k+1)} - v^{(n+1)}\|_{B_{p,r}^{s-1}}) \cdot (\|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2) \end{aligned}$$

Thus, we have that

$$\begin{aligned} & \exp\left(-C \int_0^t \left\| \frac{1}{2}(u_x^{(n+k)} v_x^{(n+k)} - u^{(n+k)} v^{(n+k)}) \right\|_{B_{p,r}^{s-1}} d\tau\right) \cdot (\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}} + \\ & \|v^{(n+k+1)} - v^{(n+1)}\|_{B_{p,r}^{s-1}})(t) \|_{B_{p,r}^{s-1}} \\ & \leq (\|u_0^{(n+k+1)} - u_0^{(n+1)}\|_{B_{p,r}^{s-1}} + \|v_0^{(n+k+1)} - v_0^{(n+1)}\|_{B_{p,r}^{s-1}}) + \\ & C \int_0^t \exp\left(-C \int_0^\tau \left\| \frac{1}{2}(u_x^{(n+k)} v_x^{(n+k)} - u^{(n+k)} v^{(n+k)}) \right\|_{B_{p,r}^{s-1}} d\tau'\right) \cdot [(\|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}} + \\ & \|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}})(\|u^{(n)}\|_{B_{p,r}^s}^2 + \|u^{(n+1)}\|_{B_{p,r}^s}^2 + \|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2 + \\ & \|v^{(n+1)}\|_{B_{p,r}^s}^2 + \|v^{(n)}\|_{B_{p,r}^s}^2 + |b|) + (\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}} + \|v^{(n+k+1)} - v^{(n+1)}\|_{B_{p,r}^{s-1}}) \times \\ & (\|u^{(n+k)}\|_{B_{p,r}^s}^2 + \|v^{(n+k)}\|_{B_{p,r}^s}^2)] d\tau. \end{aligned} \quad (24)$$

Since $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is uniformly bounded in $E_{p,r}^s(T) \times E_{p,r}^s(T)$ and

$$\begin{aligned} u_0^{(n+k+1)} - u_0^{(n+1)} &= S_{n+k+1} u_0 - S_{n+1} u_0 = \sum_{q=n+1}^{n+k} \Delta_q u_0, \\ v_0^{(n+k+1)} - v_0^{(n+1)} &= S_{n+k+1} v_0 - S_{n+1} v_0 = \sum_{q=n+1}^{n+k} \Delta_q v_0, \end{aligned}$$

there exists a constant $C(T)$ independent of n and k such that for all $t \in [0, T]$

$$\begin{aligned} & \|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}}(t) + \|v^{(n+k+1)} - v^{(n+1)}\|_{B_{p,r}^{s-1}}(t) \\ & \leq C(T) \left(2^{-n} + \int_0^t \|u^{(n+k)} - u^{(n)}\|_{B_{p,r}^{s-1}}(\tau) + \|v^{(n+k)} - v^{(n)}\|_{B_{p,r}^{s-1}}(\tau) d\tau \right). \end{aligned}$$

And by induction with regard to index n , one can deduce that

$$\begin{aligned} & \|u^{(n+k+1)} - u^{(n+1)}\|_{L_T^\infty(B_{p,r}^{s-1})} + \|v^{(n+k+1)} - v^{(n+1)}\|_{L_T^\infty(B_{p,r}^{s-1})} \\ & \leq \frac{(tC(T))^{n+1}}{(n+1)!} (\|u^{(k)}\|_{L_T^\infty(B_{p,r}^s)} + \|v^{(k)}\|_{L_T^\infty(B_{p,r}^s)}) + C(T) \sum_{i=0}^n 2^{-(n-i)} \frac{(tC(T))^i}{i!}. \end{aligned}$$

Similarly, $\|u^{(k)}\|_{L_T^\infty(B_{p,r}^s)}, \|v^{(k)}\|_{L_T^\infty(B_{p,r}^s)}$ can be bounded independent of k . Hence, we conclude that there exists some new constant $C'(T)$ independent of n and k such that

$$\|u^{(n+k+1)} - u^{(n+1)}\|_{B_{p,r}^{s-1}}(t) + \|v^{(n+k+1)} - v^{(n+1)}\|_{B_{p,r}^{s-1}}(t) \leq 2^{-n} C'(T).$$

Therefore, $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is a Cauchy sequence in $C([0, T]; B_{p,r}^{s-1})^2$.

Proof of Theorem 3.1 According to Lemma 3.3, we have that $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is a Cauchy sequence in $C([0, T]; B_{p,r}^{s-1}) \times C([0, T]; B_{p,r}^{s-1})$, so it converges to some function $(u, v) \in C([0, T]; B_{p,r}^{s-1}) \times C([0, T]; B_{p,r}^{s-1})$. The only thing that we need to do is to check that (u, v) belongs to $E_{p,r}^s(T) \times E_{p,r}^s(T)$ and solves the Cauchy problem (1). Since $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ is uniformly bounded in

$L^\infty([0, T]; B_{p,r}^s) \times L^\infty([0, T]; B_{p,r}^s)$ according to Lemma 3.3, the Fatou property for the Besov spaces guarantees that (u, v) also belongs to $L^\infty([0, T]; B_{p,r}^s) \times L^\infty([0, T]; B_{p,r}^s)$.

On the other hand, as $(u^{(n)}, v^{(n)})_{n \in \mathbb{N}}$ converges to (u, v) in $C([0, T]; B_{p,r}^{s-1}) \times C([0, T]; B_{p,r}^{s-1})$, an interpolation argument ensures that the convergence holds in $C([0, T]; B_{p,r}^{s'} \times C([0, T]; B_{p,r}^{s-1})$, for any $s' < s$. It is then easy to cross the limit in the equation (15) and to conclude that (u, v) is indeed a solution to the Cauchy problem (1). Since (u, v) belongs to $L^\infty([0, T]; B_{p,r}^s) \times L^\infty([0, T]; B_{p,r}^s)$, the right-hand side of the equation

$$\partial_t m^{(n+1)} + \frac{1}{2}(u_x^{(n)} v_x^{(n)} - u^{(n)} v^{(n)}) \partial_x m^{(n+1)} = P(u^{(n)}, v^{(n)})$$

belongs to $L^\infty([0, T]; B_{p,r}^{s-2}) \times L^\infty([0, T]; B_{p,r}^{s-2})$. Especially, for the case $r < \infty$, Lemma 2.2 implies that $(u, v) \in C([0, T]; B_{p,r}^{s'} \times C([0, T]; B_{p,r}^{s-1})$ for any $s' < s$. Finally, using the equation again, we see that $(\partial_t u, \partial_t v) \in C([0, T]; B_{p,r}^{s-1}) \times C([0, T]; B_{p,r}^{s-1})$ if $r < \infty$, and in $L^\infty([0, T]; B_{p,r}^{s-1}) \times L^\infty([0, T]; B_{p,r}^{s-1})$ otherwise. Moreover, a standard use of a sequence of viscosity approximate solutions $(u_\epsilon, v_\epsilon)_{\epsilon > 0}$ for the Cauchy problem (1) which converges uniformly in

$$C([0, T]; B_{p,r}^s) \cap C^1([0, T]; B_{p,r}^{s-1}) \times C([0, T]; B_{p,r}^s) \cap C^1([0, T]; B_{p,r}^{s-1})$$

leads to the continuity of the solution (u, v) in $E_{p,r}^s(T) \times E_{p,r}^s(T)$. The proof is completed. \square

Acknowledgements We thank the referees for their time and comments.

References

- [1] Changzheng QU, Junfeng SONG, Ruoxia YAO. *Multi-component integrable systems and invariant curve flows in certain geometries*. SIGMA Symmetry Integrability Geom. Methods Appl., 2013, **9**: 1–19.
- [2] B. FUCHSSTEINER, A. S. FOKAS. *Symplectic structures, their Bäcklund transformations and hereditary symmetries*. Phys. D., 1981/82, **4**: 47–66.
- [3] R. CAMASSA, D. HOLM. *An integrable shallow water equation with peaked solitons*. Phys. Rev. Lett., 1993, **71**(11): 1661–1664.
- [4] Huihui DAI. *Model equations for nonlinear dispersive waves in a compressible Mooney-Rivlin rod*. Acta Mech., 1998, **127**(1-4): 193–207.
- [5] A. CONSTANTIN. *On the scattering problem for the Camassa-Holm equation*. R. Soc. Lond. Proc. Ser. A Math. Phys. Eng. Sci., 2001, **457**(2008): 953–970.
- [6] A. CONSTANTIN. *The Hamiltonian structure of the Camassa-Holm equation*. Exposition. Math., 1997, **15**(1): 53–85.
- [7] R. CAMASSA, D. D. HOLM, J. M. HYMAN. *A new integrable shallow water equation*. Adv. Appl. Mech., 1994, **31**: 1.
- [8] H. P. MCKEAN. *Fredholm determinants and the Camassa-Holm hierarchy*, Comm. Pure Appl. Math., 2003, **56**(5): 638–680.
- [9] A. CONSTANTIN. *Existence of permanent and breaking waves for a shallow water equation: a geometric approach*. Ann. Inst. Fourier (Grenoble), 2000, **50**(2): 321–362.
- [10] YI. A. LI, P. J. OLIVER. *Well-posedness and blow-up solutions for an integrable nonlinearly dispersive model wave equation*. J. Differential Equations, 2000, **162**(1): 27–63.
- [11] A. BRESSAN, A. CONSTANTIN. *Global conservative solutions of the Camassa-Holm equation*. Arch. Ration. Mech. Anal., 2007, **183**(2): 215–239.
- [12] B. Fuchssteiner, *Some tricks from the symmetry-toolbox for nonlinear equations: generalizations of the Camassa-Holm equation*. Phys. D, 1996, **95**(3-4): 229–243.
- [13] P. J. Olver and P. Rosenau, *Tri-Hamiltonian duality between solitons and solitary-wave solutions having compact support*. Phys. Rev. E (3), 1996, **53**(2): 1900–1906.
- [14] A. S. FOKAS. *On a class of physically important integrable equations*. Phys. D, 1995, **87**(1-4): 145–150.

- [15] Zhijun QIAO. *A new integrable equation with cuspons and W/M-shape-peaks solitons*. J. Math. Phys., 2006, **47**(11): 1–9.
- [16] Zhijun QIAO, Xianqi LI. *An integrable equation with nonsmooth solitons*. Theor. Math. Phys., 2011, **167**: 584–589.
- [17] Ying FU, Guilong GUI, Yue LIU, et al, *On the Cauchy problem for the integrable modified Camassa-Holm equation with cubic nonlinearity*. J. Differential Equations, 2013, **255**(7): 1905–1938.
- [18] Guilong GUI, Yue LIU, P. J. OLIVER, et al. *Wave-breaking and peakons for a modified Camassa-Holm equation*. Comm. Math. Phys., 2013, **319**(3): 731–759.
- [19] Zhijun QIAO, Baoqiang XIA, Jibin LI. *Integrable system with peakon, weak kink, and kink-peakon interactional solutions*. Front.math.china, 2013, **8**(5): 1185–1196.
- [20] A. CONSTANTIN, R. I. IVANOV. *On an integrable two-component Camassa-Holm shallow water system*. Phys. Lett. A, 2008, **372**(48): 7129–7132.
- [21] J. ESCHER, O. LECHTENFELD, Zhaoyang YIN. *Well-posedness and blow-up phenomena for the 2-component Camassa-Holm equation*. Discrete Contin. Dyn. Syst., 2007, **19**(3): 493–513.
- [22] R. IVANOV. *Two-component integrable system modelling shallow water waves: the constant vorticity case*. Wave Motion, 2009, **46**(6): 389–396.
- [23] Guilong GUI, Yue LIU. *On the Cauchy problem for the two-component Camassa-Holm system*. Math. Z., 2011, **268**(1-2): 45–66.
- [24] R. DANCHIN. *A few remarks on the Camassa-Holm equation*. Differential Integral Equations, 2001, **14**(8): 953–988.
- [25] H. BAHOURI, J. Y. CHEMIN, R. DANCHIN. *Fourier Analysis and Nonlinear Partial Differential Equations*. Springer, Heidelberg, 2011.
- [26] R. DANCHIN. *Fourier analysis methods for PDEs*. Lecture Notes, 14 November, 2005.